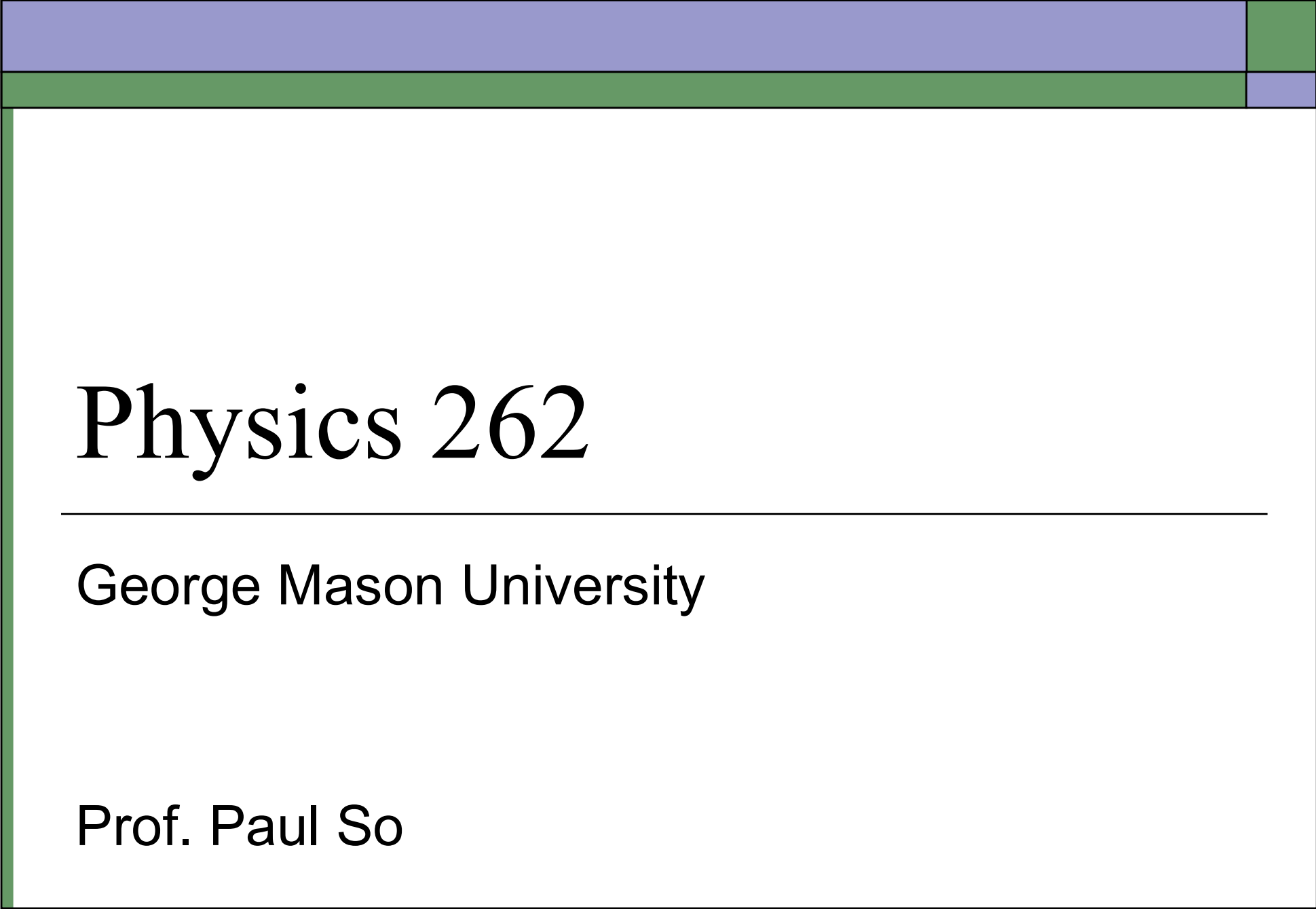




Physics 262

George Mason University

Prof. Paul So

Chapter 35: Interference

- ❑ Interference and Coherent Sources
- ❑ Two-Source Interference of Light
- ❑ Intensity of Interference Patterns
- ❑ Interference in Thin Films
- ❑ The Michelson Interferometer

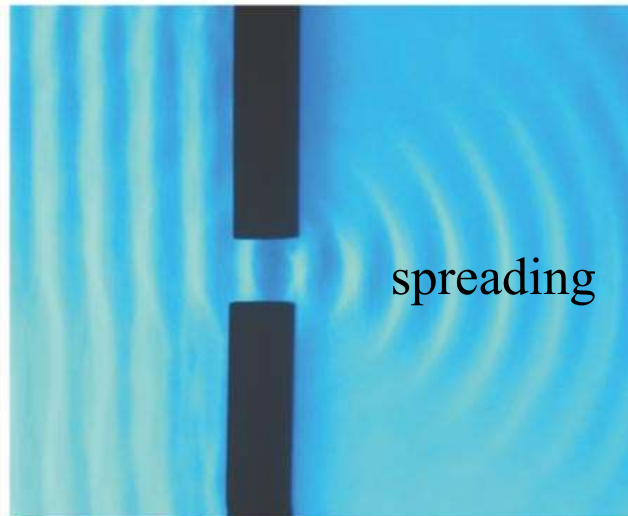


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Wave Nature of Light

- Previous Chapters (Geometric Optics) $\lambda \ll L$
 - Rays Model is an approximation of EM waves with rays pointing in the direction of propagation
- Next Couple of Chapters (Wave/Physical Optics) $\lambda \sim L$
 - Like water waves, light *spreads* and *interferes* with each other.
 - Observed phenomena *cannot* be accounted for by rays:

Diffraction



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Interference

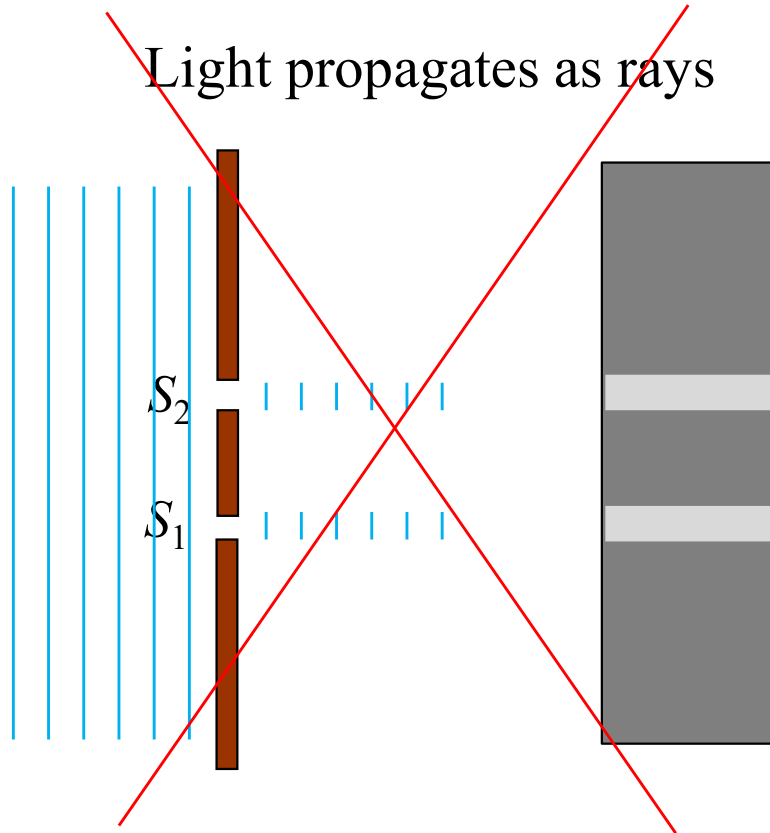


constructive/
destructive
interference
patterns

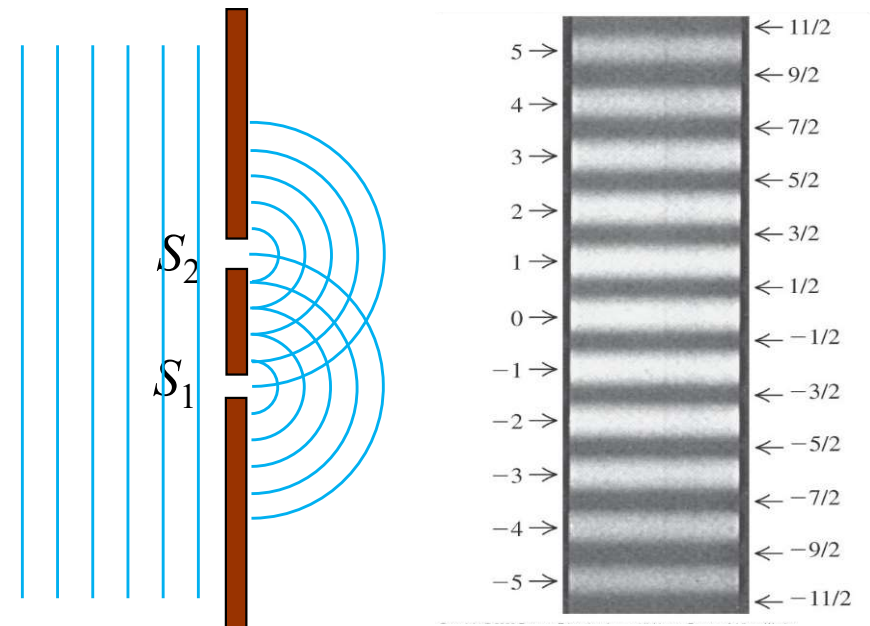
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Wave Nature of Light: Diffraction & Interference

Light propagates as rays

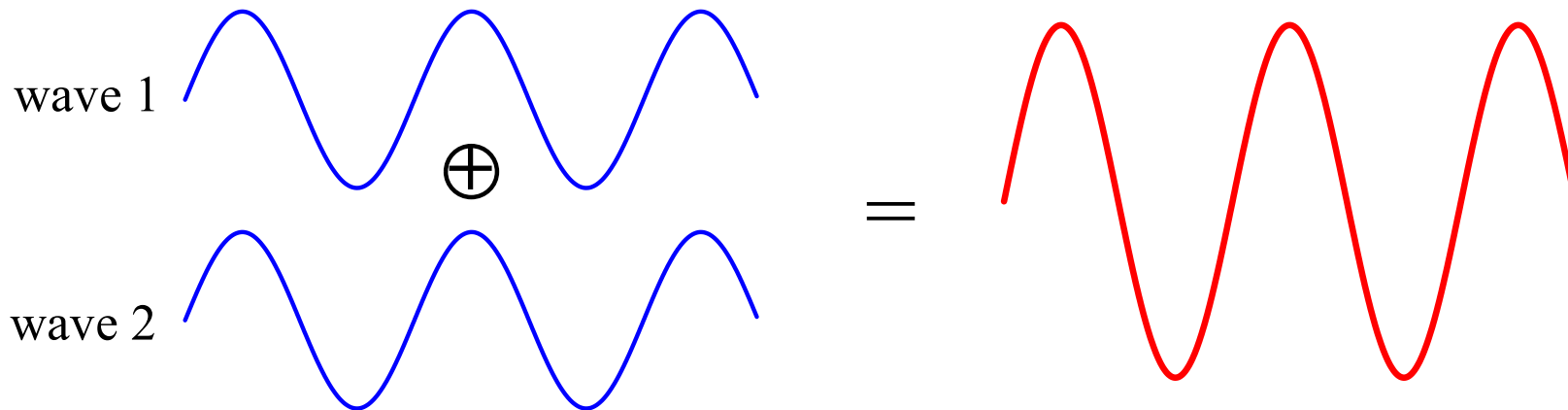


Light propagates as waves

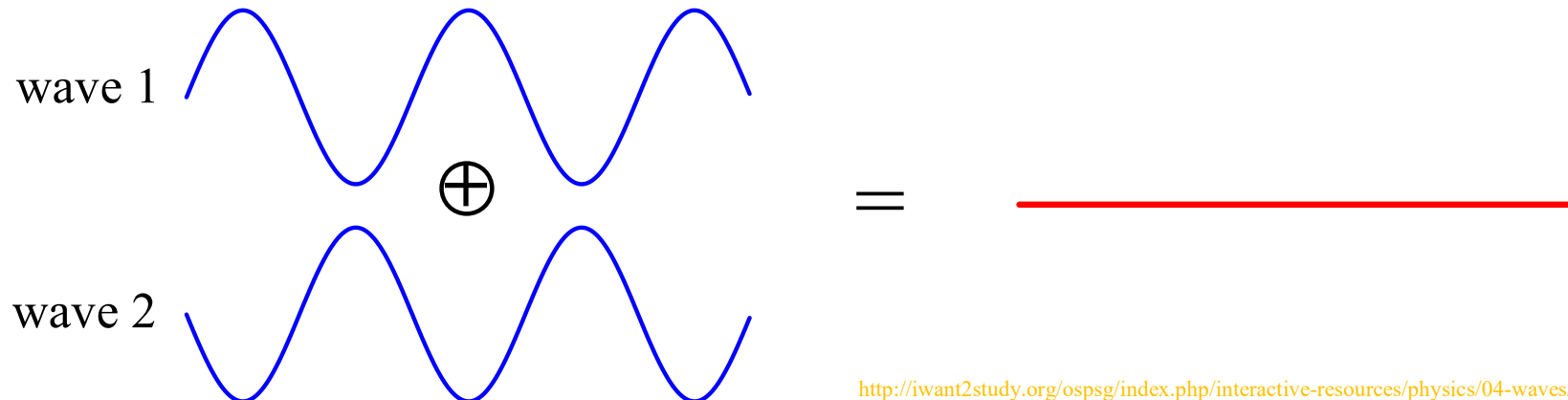


Interference and Superposition

Constructive Interference (+ peaks aligns w/ + peaks)

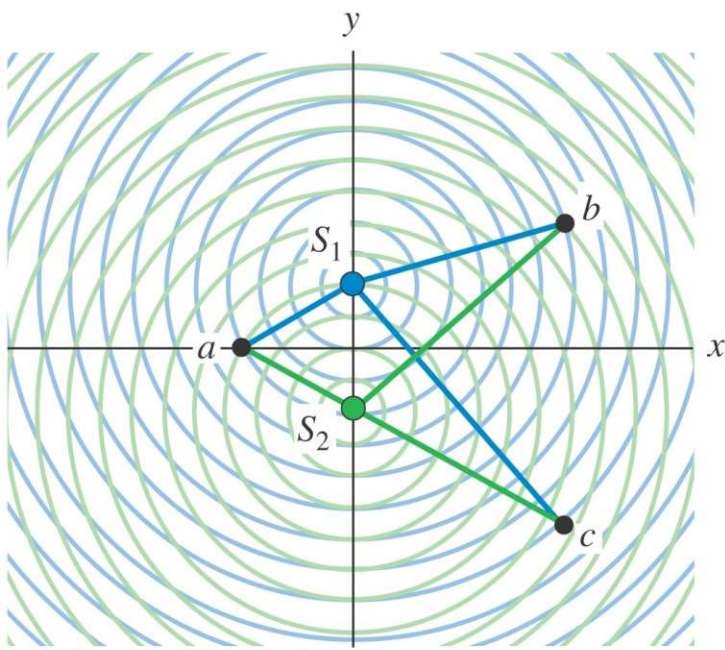


Destructive Interference (+ peaks aligns w/ - peaks: $\lambda/2$ apart)



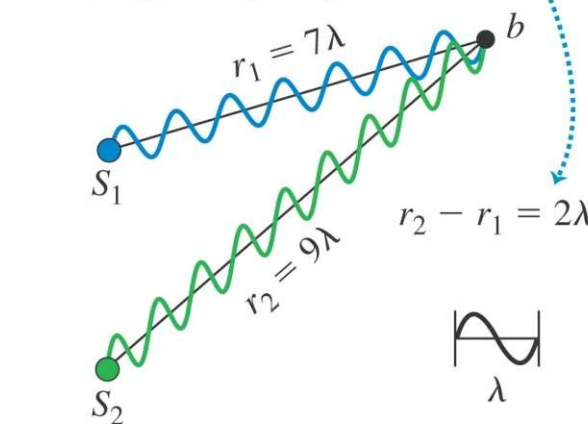
Interference, Superposition, & Path Difference

D two sheets



(a) Point a is symmetric with respect to the two coherent sources. Waves will arrive in phase *constructively*: $r_2 - r_1 = 0$.

(b) Conditions for constructive interference: Waves interfere constructively if their path lengths differ by an integral number of wavelengths: $r_2 - r_1 = m\lambda$.

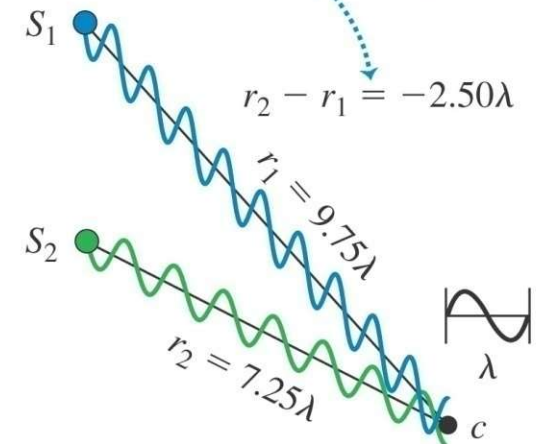


Constructive Inter.

$$r_2 - r_1 = m\lambda$$

$$(m = 0, \pm 1, \pm 2, \dots)$$

(c) Conditions for destructive interference: Waves interfere destructively if their path lengths differ by a half-integral number of wavelengths: $r_2 - r_1 = (m + \frac{1}{2})\lambda$.



Destructive Inter.

$$r_2 - r_1 = \left(m + \frac{1}{2}\right)\lambda$$

$$(m = 0, \pm 1, \pm 2, \dots)$$

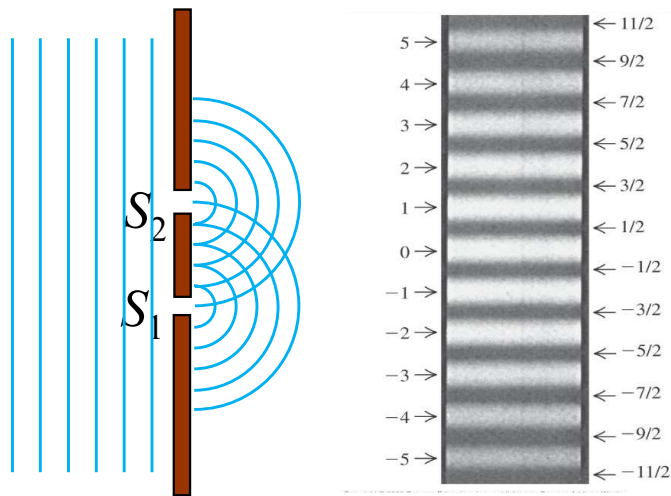
r_1 : distance to S_1

r_2 : distance to S_2

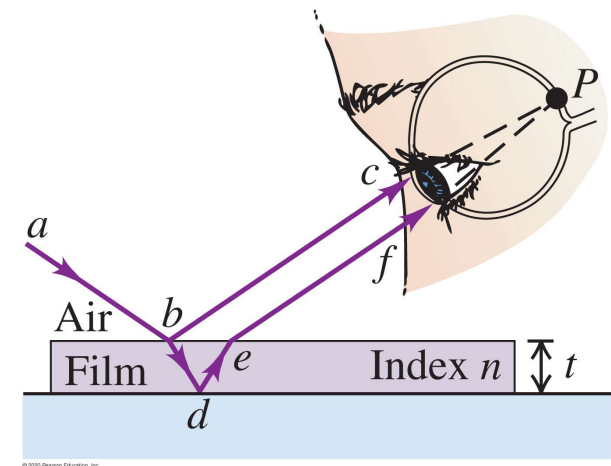
$r_2 - r_1 =$ **path difference**

Many Ways when Waves can Interfere

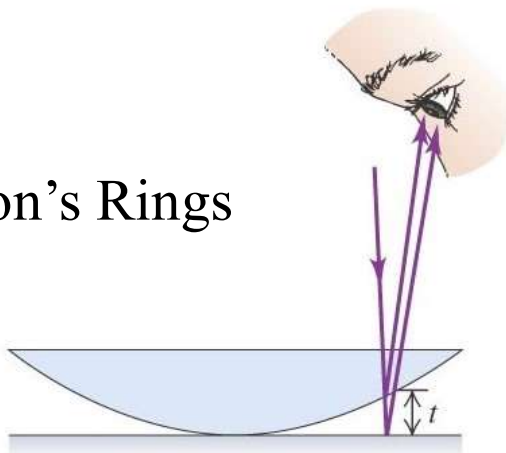
Young's Double Slits



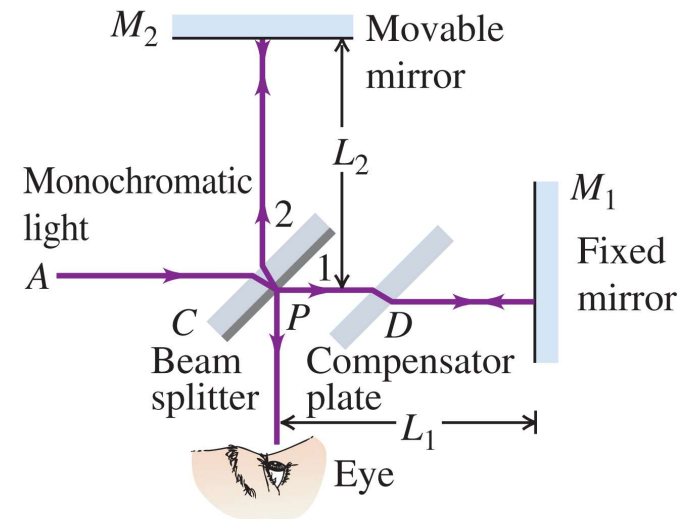
Diff Thin Film Scenarios



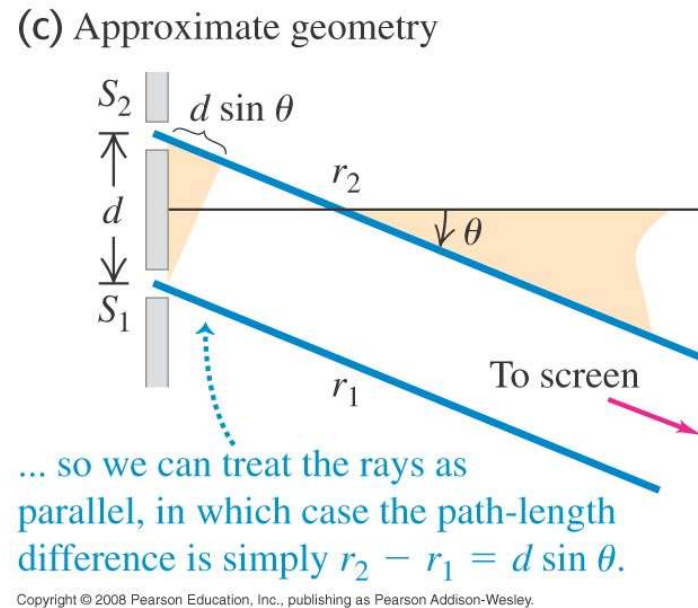
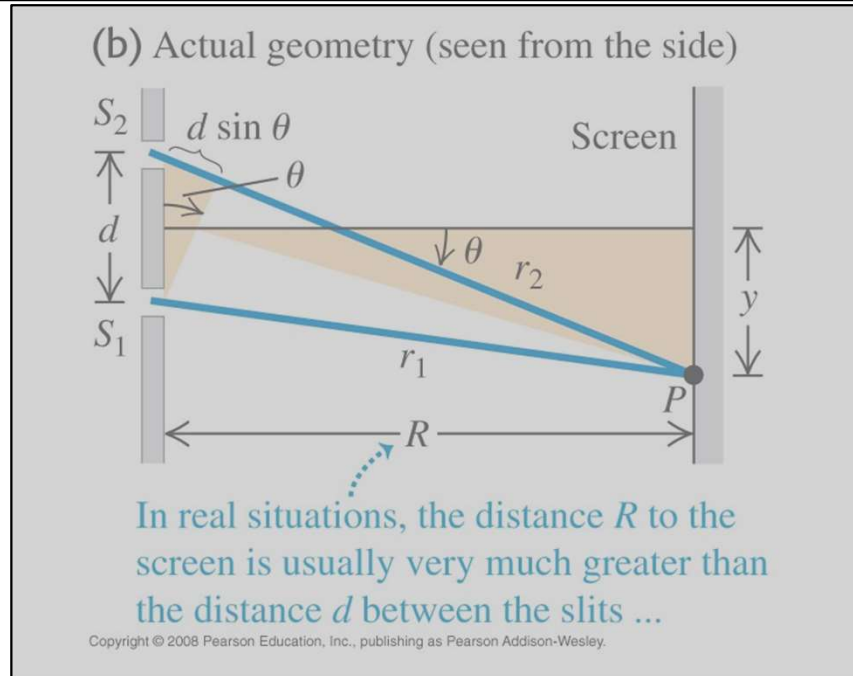
Newton's Rings



Michelson Interferometer



Young's Double-Slit



If screen is far away so that $R \gg d$, we can assume rays from S_1 and S_2 to be approximately parallel.

Then, from the simplified geometry (right panel), we have an explicit expression for the **path difference**:

$$r_2 - r_1 = d \sin \theta$$

(θ is the angular location of observation point P on the screen.)

Constructive/Destructive Young's Double-Slit Interference

For observing distance $R \gg d$, we have the following conditions:

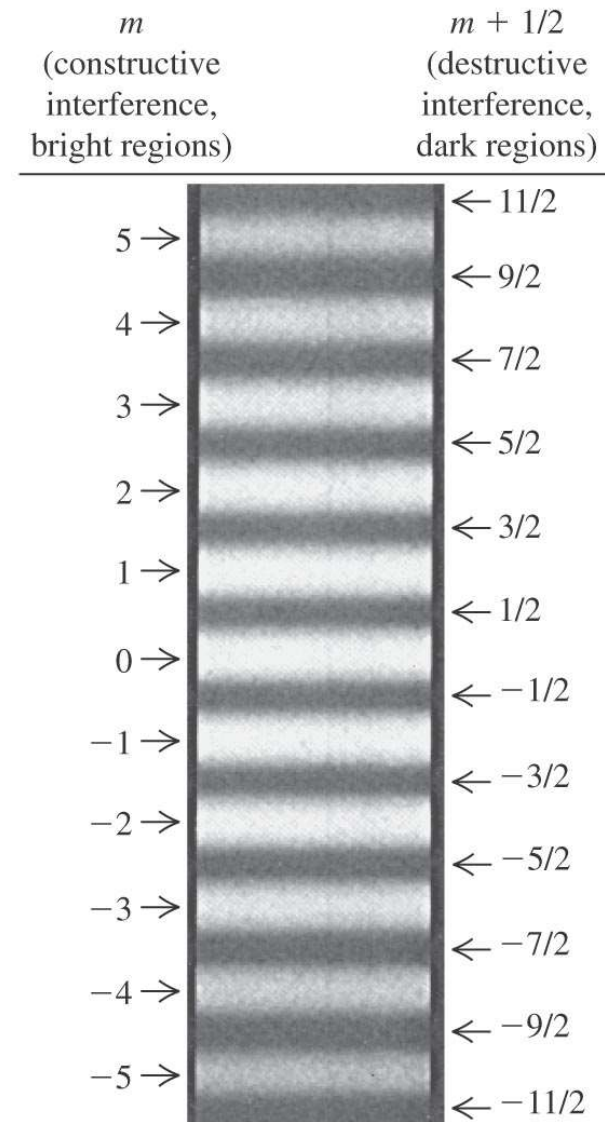
Constructive Interference: Two Slit Interference

$$d \sin \theta = m\lambda \quad (m = 0, \pm 1, \pm 2, \dots)$$

Destructive Interference: Two Slit Interference

$$d \sin \theta = \left(m + \frac{1}{2}\right)\lambda \quad (m = 0, \pm 1, \pm 2, \dots)$$

- The bright/dark bands in the pattern are called fringes
- m is the *order* of the fringes



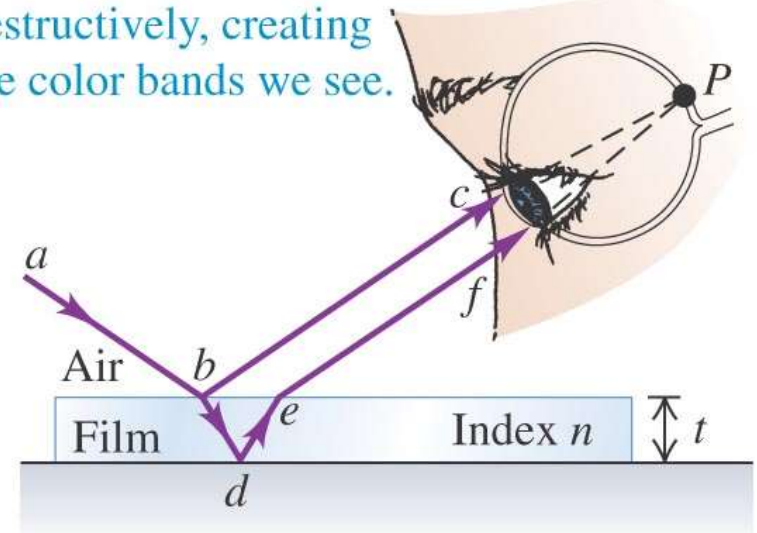
Interference in Thin Films

Color fringes observed from an oil slick on water or on a soap bubble are the white-light *interference* patterns produced by the *reflected* light off a *thin film* of oil or soap.



Light reflected from the upper and lower surfaces of the film comes together in the eye at P and undergoes interference.

Some colors interfere constructively and others destructively, creating the color bands we see.



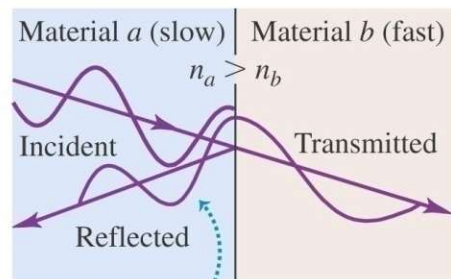
Phase Reversal upon Reflection

From Maxwell's Equations, one can show that the reflected wave will suffer a 180° or $\lambda/2$ phase shift if it is reflected off from a medium with a *larger* n .

$$E_r = \frac{n_a - n_b}{n_a + n_b} E_i \quad (\text{for normal incidence})$$

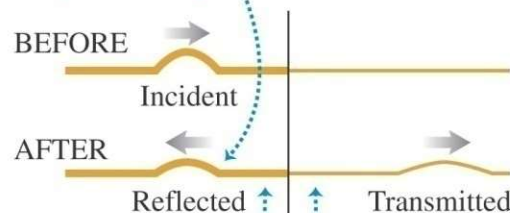
Electromagnetic waves propagating in optical materials

(a) If the transmitted wave moves *faster* than the incident wave ...



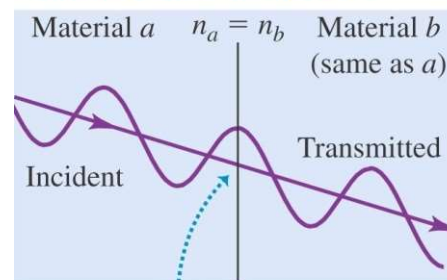
... the reflected wave undergoes no phase change.

Mechanical waves propagating on ropes

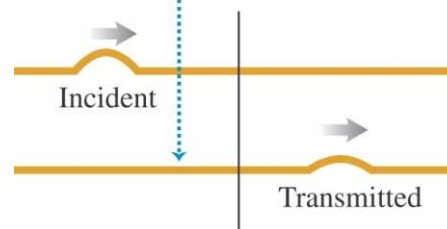


Waves travel slower on thick ropes than on thin ropes.

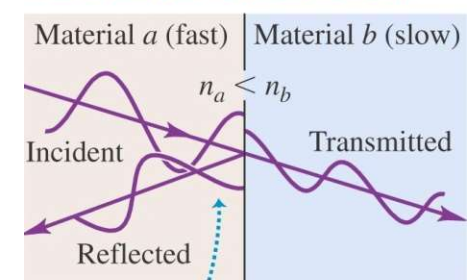
(b) If the incident and transmitted waves have the same speed ...



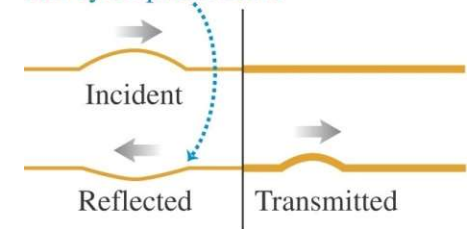
... there is no reflection.



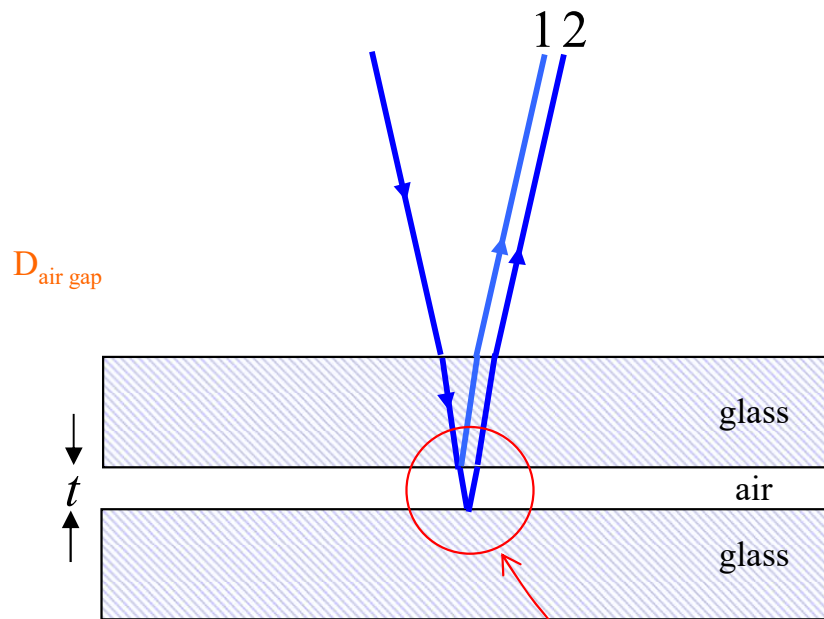
(c) If the transmitted wave moves *slower* than the incident wave ...



... the reflected wave undergoes a half-cycle phase shift.



Interference from a Thin Air Gap



Assumptions:

- Thickness of air gap t is small
- Thickness of glass is large
- Incident light is nearly *normal* at the upper plate.

Interference due to this small gap

NOTE: While wave #2 suffers a phase reversal upon reflection at the bottom of the gap, wave #1 did not.

Interference from a Thin Air Gap

Condition for Constructive and Destructive from thin film when **one** of the waves suffers a phase reversal upon reflection

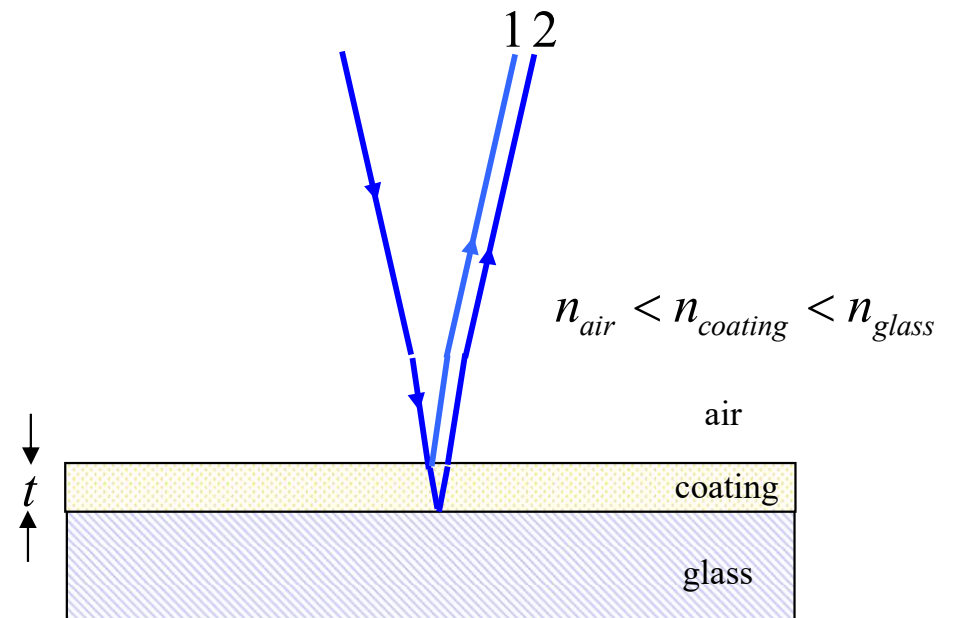
Constructive: $2t = \left(m + \frac{1}{2}\right)\lambda, \quad m = 0, 1, 2, \dots$

Note that we need to take the extra “sign change” or half a wave ($\lambda/2$) into account on the right-hand side

Destructive: $2t = m\lambda, \quad m = 0, 1, 2, \dots$

→ the switched condition on the right-hand side !

Another Thin Film Example (non-refractive coating on lens)



wave #1: reflected from top interface of the coating:

$$n_{air} < n_{coating} \Rightarrow 180^\circ \text{ (or } \pi \text{) phase reversal}$$

wave #2: reflected from bottom interface of the coating:

$$n_{coating} < n_{glass} \Rightarrow 180^\circ \text{ (or } \pi \text{) phase reversal}$$

Interference from a Thin Film

Condition for Constructive and Destructive from thin film when **both** waves suffers a phase reversal upon reflection (NO net sign change)

Constructive:

$$2t = m\lambda_n, \quad m = 0, 1, 2, \dots$$
$$2n_{\text{coating}}t = m\lambda, \quad m = 0, 1, 2, \dots$$

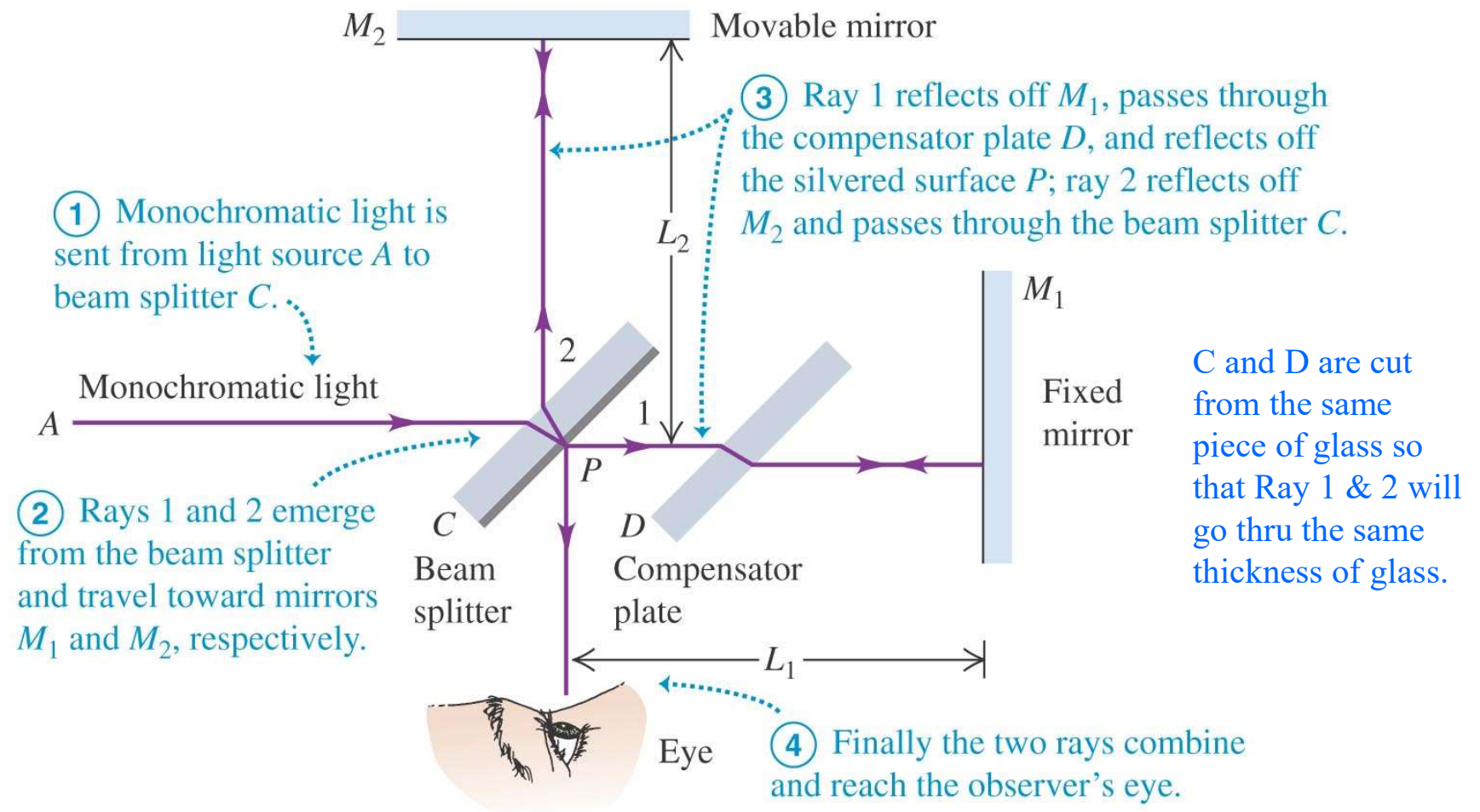
Since reflection occurs inside a medium with n , wavelength is shorter

Destructive:

$$2t = \left(m + \frac{1}{2}\right)\lambda_n, \quad m = 0, 1, 2, \dots$$
$$2n_{\text{coating}}t = \left(m + \frac{1}{2}\right)\lambda, \quad m = 0, 1, 2, \dots$$

$$\lambda_n = \frac{\lambda}{n}$$

Michelson Interferometer



Distances comparable to λ can be measured with ease using this device by counting fringes.