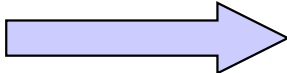# Typical Usage for the Ideal Gas Law

- For a fixed amount of gas (nR=const)

$$\frac{PV}{T}=nR=const$$

- So, if we have a gas at two different states 1(before) and 2 (after), their state variables are related simply by:

$$\frac{P_1 V_1}{T_1}=\frac{P_2 V_2}{T_2}$$

We can use this relation to solve for any unknown variables with the others being given.

Example 18.2

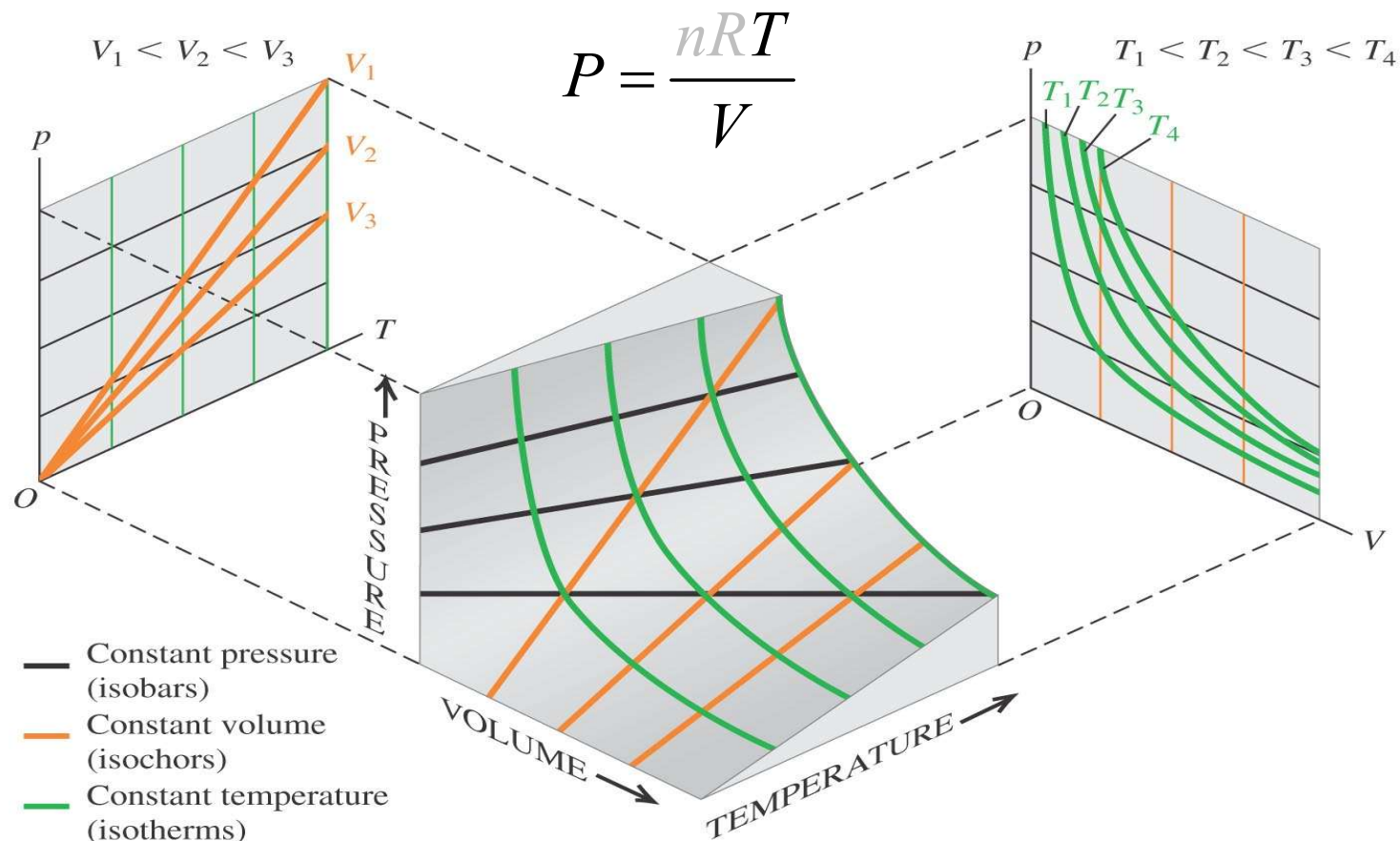In an automobile engine, a mixture of air/gasoline is being compressed before ignition.

☐ Typical compression ration 1 to 9

☐ Initial $P = 1$ atm and $T = 27$ °C

Find the temperature of the compressed gas if we are given the pressure after compression to be 21.7atm.
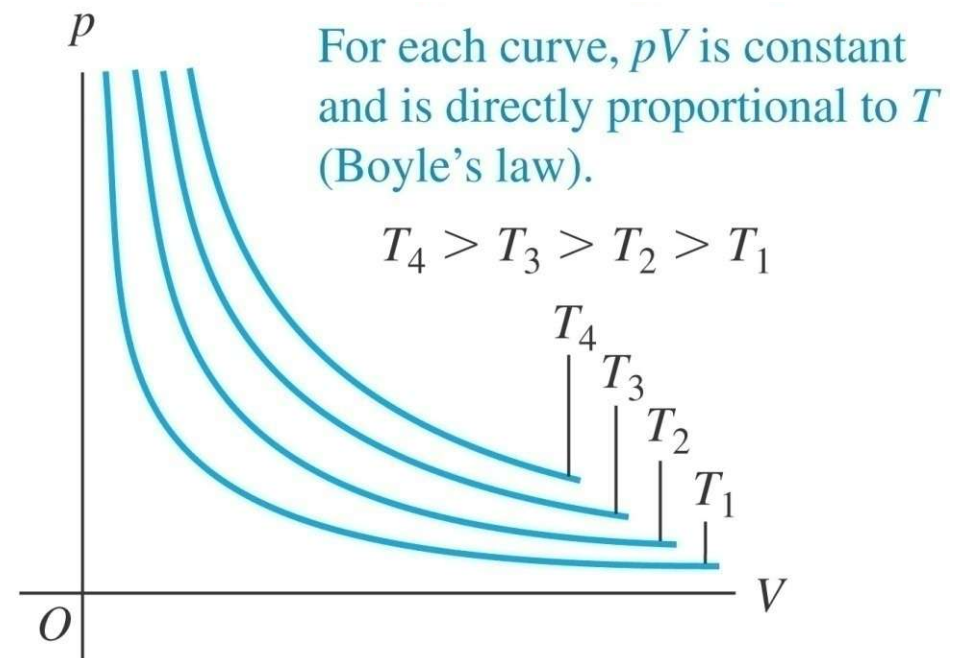
http://complex.gmu.edu/www-phys/phys262/soln/ex18.2.pdf

The Ideal Gas Law (graphical view)

P,V,T relationship in the Ideal Gas Law can be visualize graphically as a surface in 3D.

$$P = \frac{nRT}{V}$$

Constant pressure (isobars)

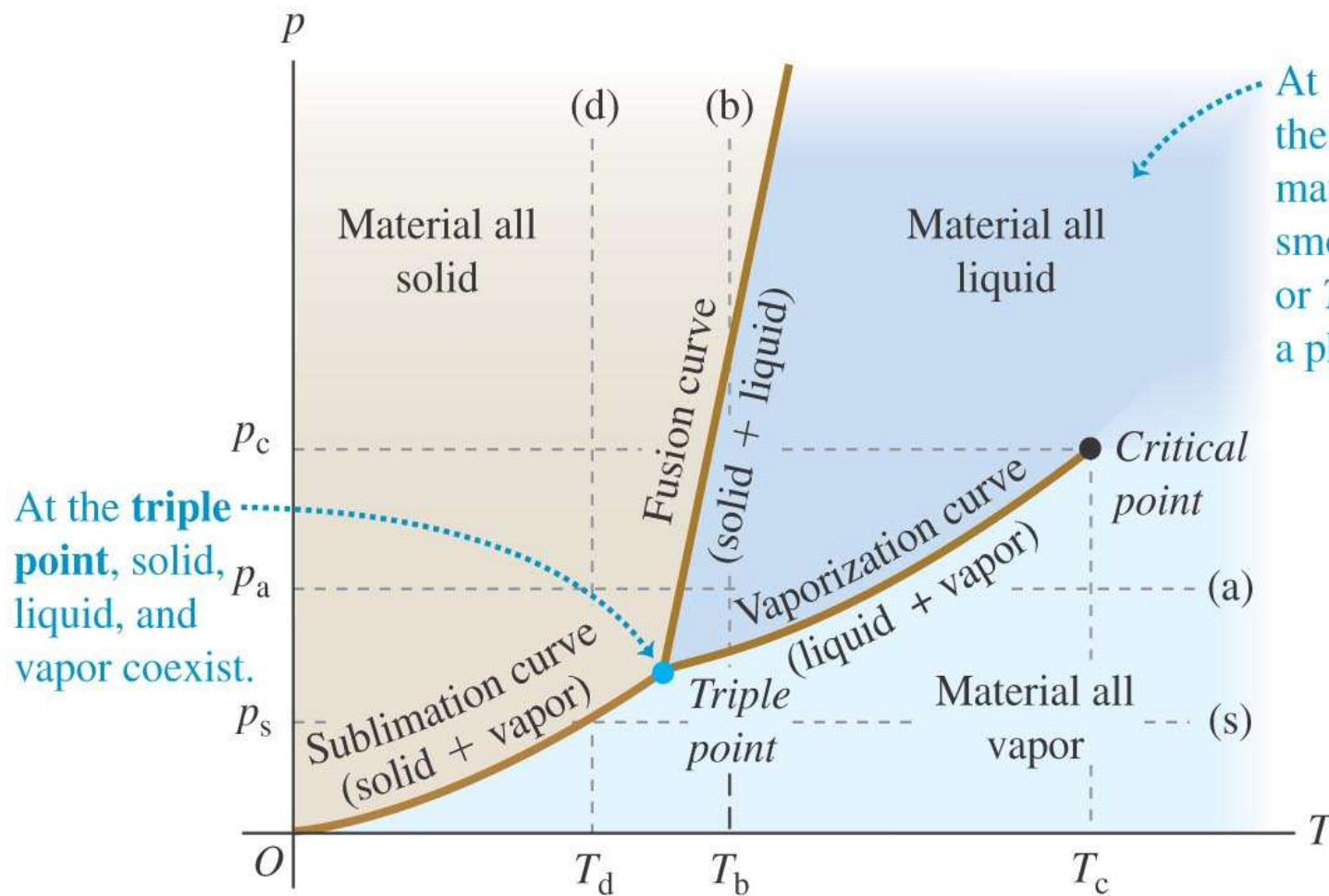Constant volume (isochors)

Constant temperature (isotherms)

PV Diagrams

□ 2D projections of the previous 3D surface.

□ Evolution of a gas at *constant T* will move along these curves called *isotherms*.

p

For each curve, pV is constant and is directly proportional to T (Boyle's law).

$T_4 > T_3 > T_2 > T_1$

T_4
T_3
T_2
T_1

O

V

Gives P vs. V at a various T: $P = (nRT)\dfrac{1}{V}$
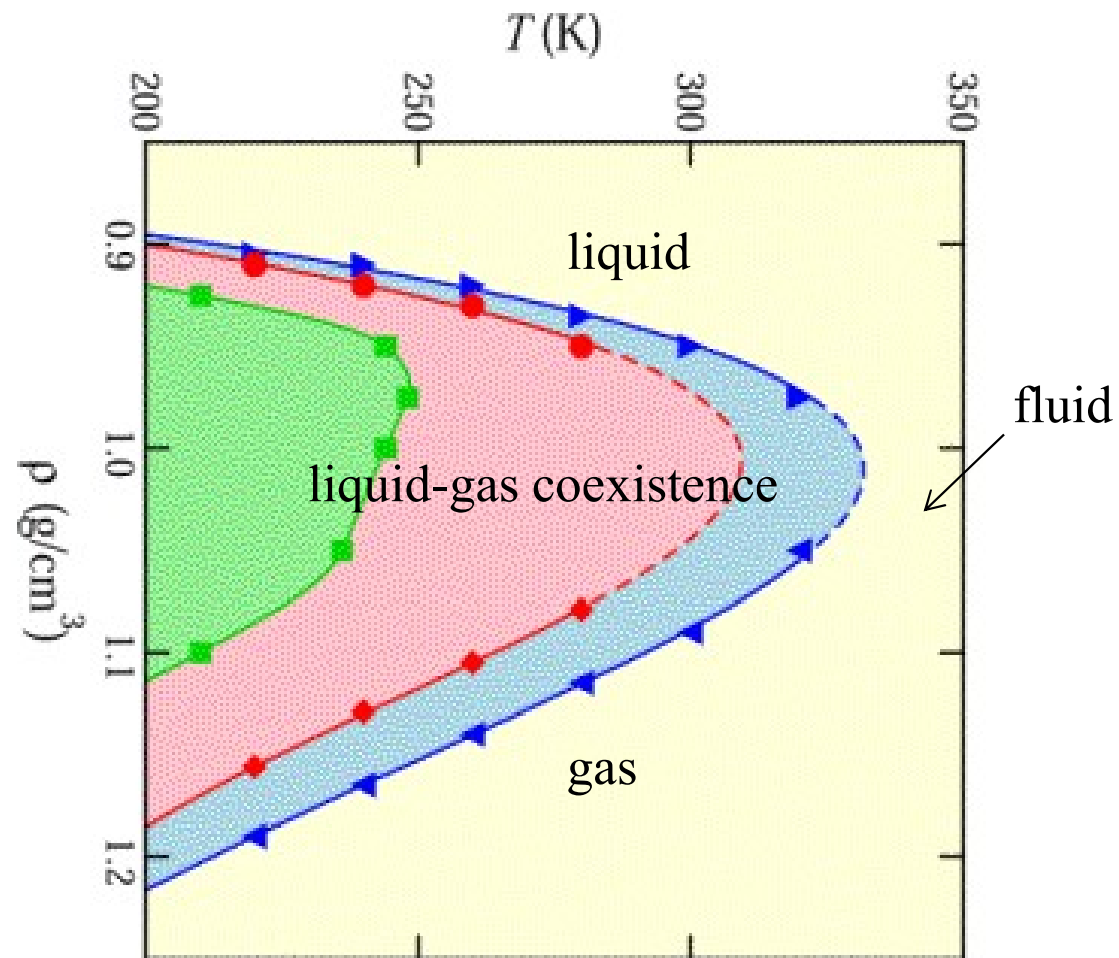
Phases of Matter (reading phase diagrams)

At the **triple point**, solid, liquid, and vapor coexist.

At T and p values above the **critical point**, the material properties change smoothly with changing p or T, rather than undergoing a phase change.
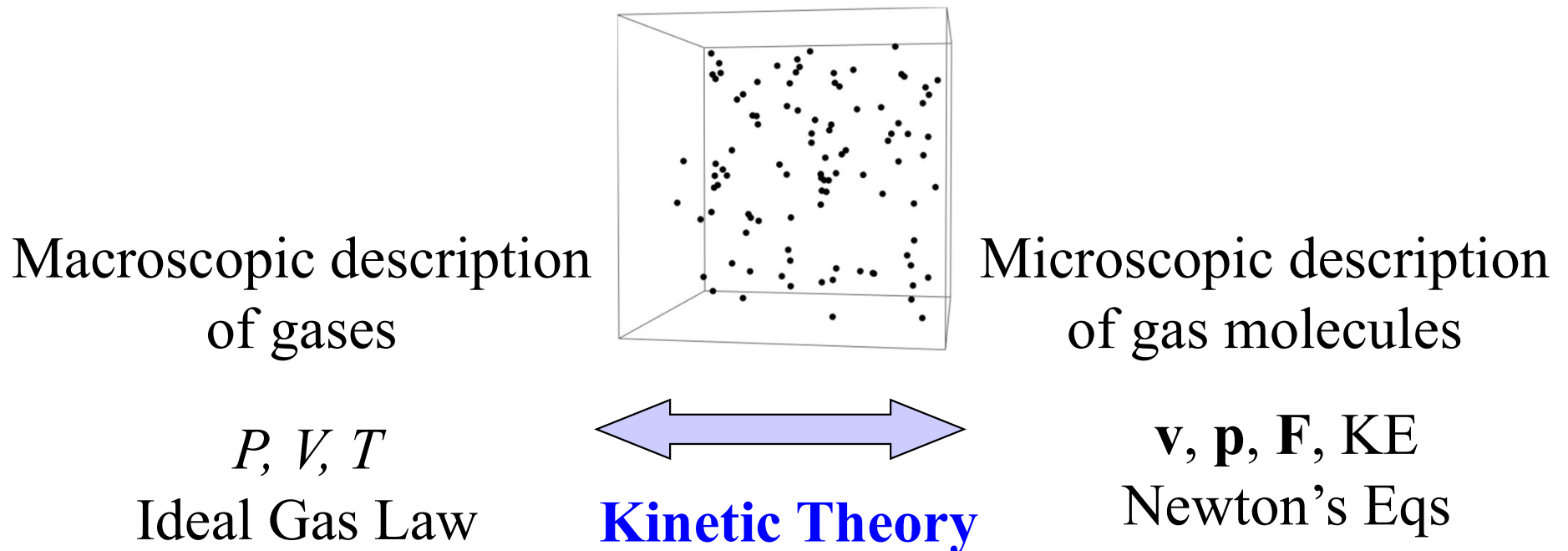
Material all solid

Material all liquid

Material all vapor

Fusion curve (solid + liquid)

Sublimation curve (solid + vapor)

Vaporization curve (liquid + vapor)

Critical point

Triple point

Critical Point and Liquid-Gas Coexistence

Kinetic-Molecular Model of an Ideal Gas

Macroscopic description
of gases

Microscopic description
of gas molecules

P, V, T

Ideal Gas Law

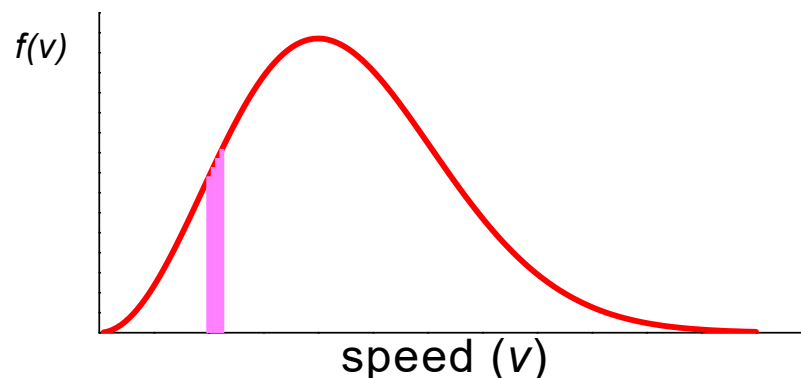**Kinetic Theory**

v, **p**, **F**, KE
Newton's Eqs

☐ An example of a successful theoretical linkage between the "micro" and "macro" descriptions for an ideal gas.

☐ Explicit expressions of P & T in terms of microscopic quantities!

Kinetic Theory (assumptions)

- A very *large # N* of *identical* molecules each with mass *m* in a container with volume *V*

- Molecules behaves as *point particles*:
  - Molecule sizes << avg separate bet particles & dim of container

- Molecules moves according to Newton's laws and they move randomly with equal probability in all directions and with a fixed **distribution** *f(v)* (histogram) of speed *v*
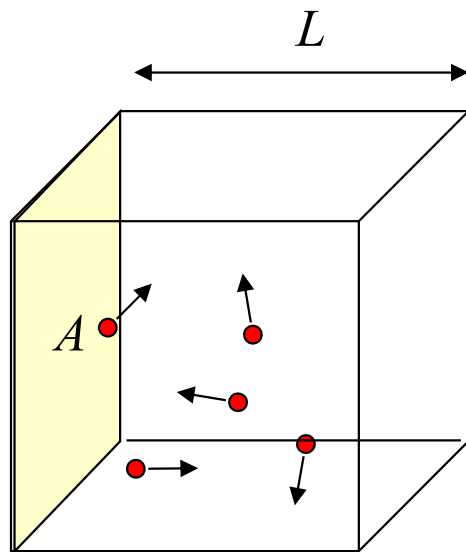
  "Fraction of molecules moving in a given range of speeds stays the same"

f(v)

speed (v)

- Molecules interact only with the walls thru elastic collisions and the container walls are perfectly rigid and infinitely massive
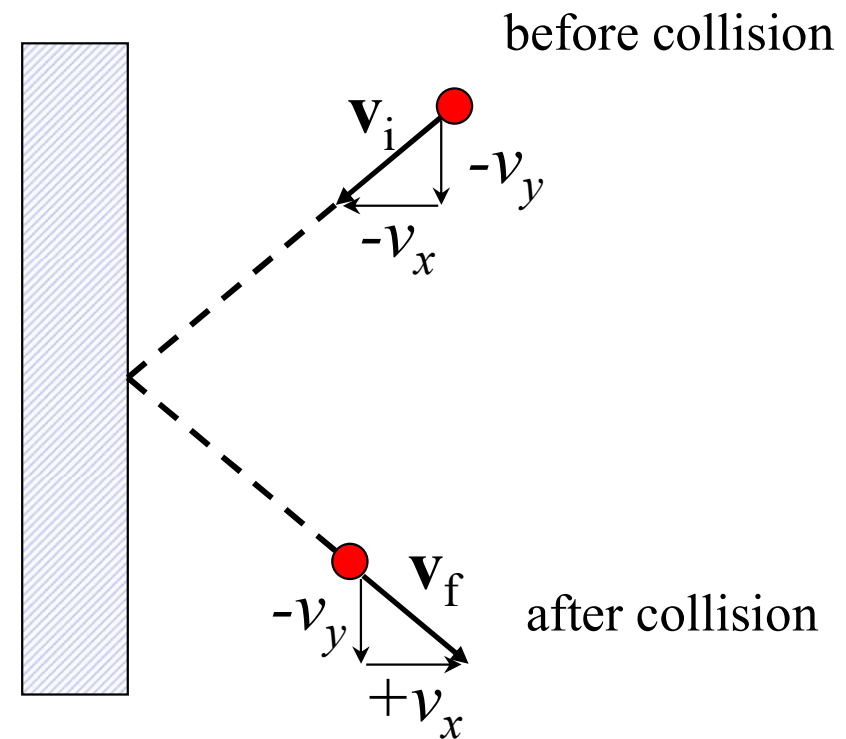 - → both KE and momentum are conserved

Kinetic Theory (model)

Idea Gas in a box with $V=AL$

Left Wall

before collision

$\mathbf{v}_i$

$-v_y$

$-v_x$

$-v_y$

$\mathbf{v}_f$

after collision

$+v_x$

- Velocity component parallel to the wall (y-component) does not change.
- Velocity component perpendicular to the wall (x-component) reverses direction.
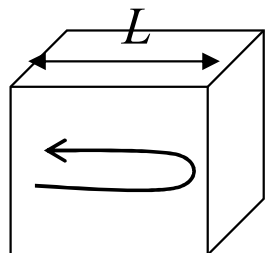- Speed v does not change.

Pressure Exerted by an Ideal Gas

"Pressure on the left wall due to molecular collisions"

Left Wall

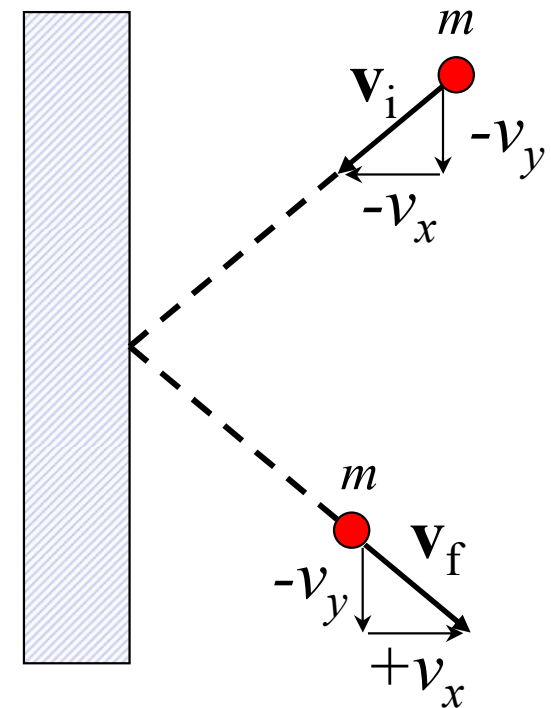1. Momentum change in *x*-dir by a molecule moving to the left at $\mathbf{v}_i$:

$$\Delta(m\mathbf{v}) = P_f - P_i = mv_x - (-mv_x) = 2mv_x$$

2. Duration, Δt, that this molecule takes (on average) to collide with the left wall again (diluted gas),

$$\Delta t = \frac{2L}{v_x}$$

Pressure Exerted by an Ideal Gas

3. Force exerted by this molecule on the left wall:

$$F = \frac{\Delta(m\mathbf{v})}{\Delta t} = \frac{2mv_x}{2L/v_x} = \frac{mv_x^2}{L}$$

4. With N molecules, total force on wall in Δt:

$$F_{tot} = \sum_{i=1}^{N} F_i = \frac{m}{L}\left(v_{1x}^2 + v_{2x}^2 + \cdots + v_{Nx}^2\right)$$

$$= \frac{m}{L}N\left(\frac{1}{N}\sum_{i=1}^{N} v_{ix}^2\right) = \frac{mN}{L}\left(v_x^2\right)_{av} \qquad Note: \quad \left(v_x^2\right)_{av} \neq \left(v_x\right)_{av}^2$$

→ Fixed *Invariant* distribution implies $(\cdot)_{av}$ will be the same when experiment is repeated.

Pressure of an Ideal Gas

5. Random direction (isotropic) assumption:

$$(v_x{}^2)_{av} = (v_y{}^2)_{av} = (v_z{}^2)_{av} \qquad (x, y, z \text{ are the same})$$

Since $v^2 = \|\vec{v}\|^2 = v_x^2 + v_y^2 + v_z^2$, we have $(v^2)_{av} = 3(v_x{}^2)_{av}$

This gives, $\quad F_{tot} = \dfrac{mN}{L} \dfrac{(v^2)_{av}}{3}$

6. Finally, the pressure on the wall is:

$$P = \frac{F_{tot}}{A} = \frac{mN}{AL} \frac{(v^2)_{av}}{3} \quad = \quad \frac{1}{3} \frac{Nm(v^2)_{av}}{V}$$

Pressure of an Ideal Gas

Rewriting, we have

(avg trans. KE per molecule)

$$PV = \frac{2}{3} N \left(\frac{1}{2} m \left(v^2 \right)_{av} \right) = \frac{2}{3} N \left(KE \right)_{av}$$

☐ This tells us that P inside a container with a fixed V:

- ■ is proportional to the # of molecules N
- ■ is proportional to the **avg.** translational KE of molecules

(These are microscopic properties of the gas.)

Molecular Interpretation of Temperature

From before: $\boxed{PV = \frac{2}{3}N\left(KE\right)_{av}}\left(= \frac{2}{3}K_{tr}\right)$

From Idea Gas Law: $PV = NkT$ (Recall k – Boltzmann Constant)

For both of these to be true, we need to have: $\frac{2}{3}\cancel{N}\left(KE\right)_{av} = \cancel{N}kT$

$\boxed{\left(KE\right)_{av} = \frac{3}{2}kT}$ (per molecule)

$K_{tr} = \frac{3}{2}NkT$ (all molecules in V)

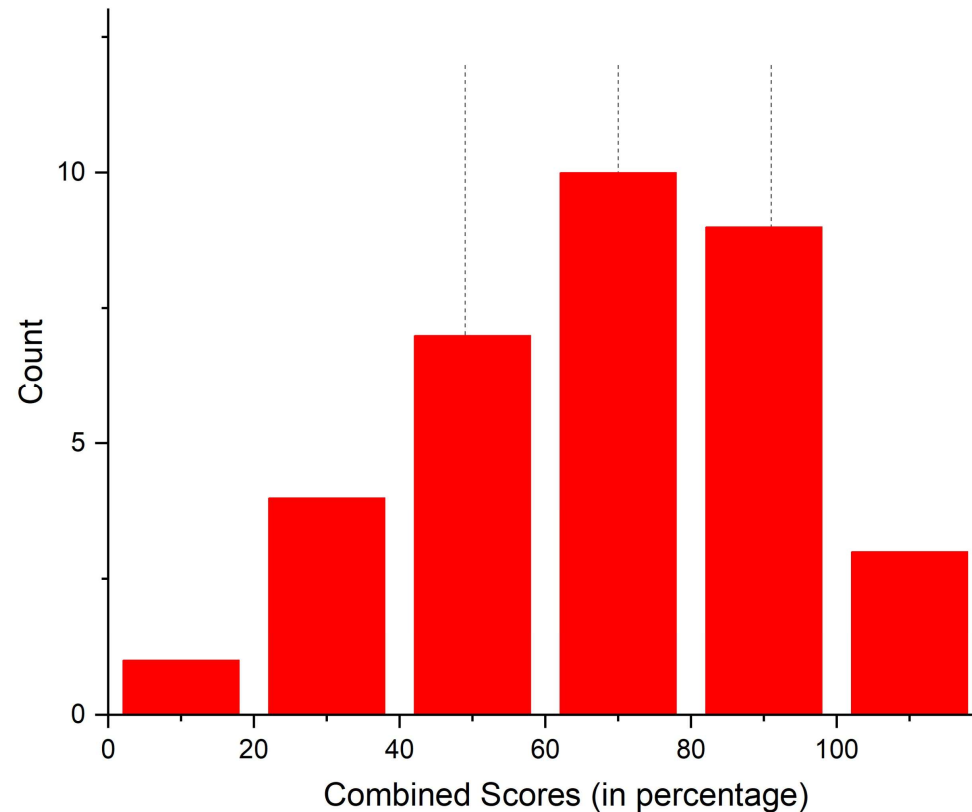(Note: K_{tr} in your book is the avg *total* KE for **all** molecules in volume V.)

"Temperature is a direct measure of the **average** translational KE of the molecules in an ideal gas."
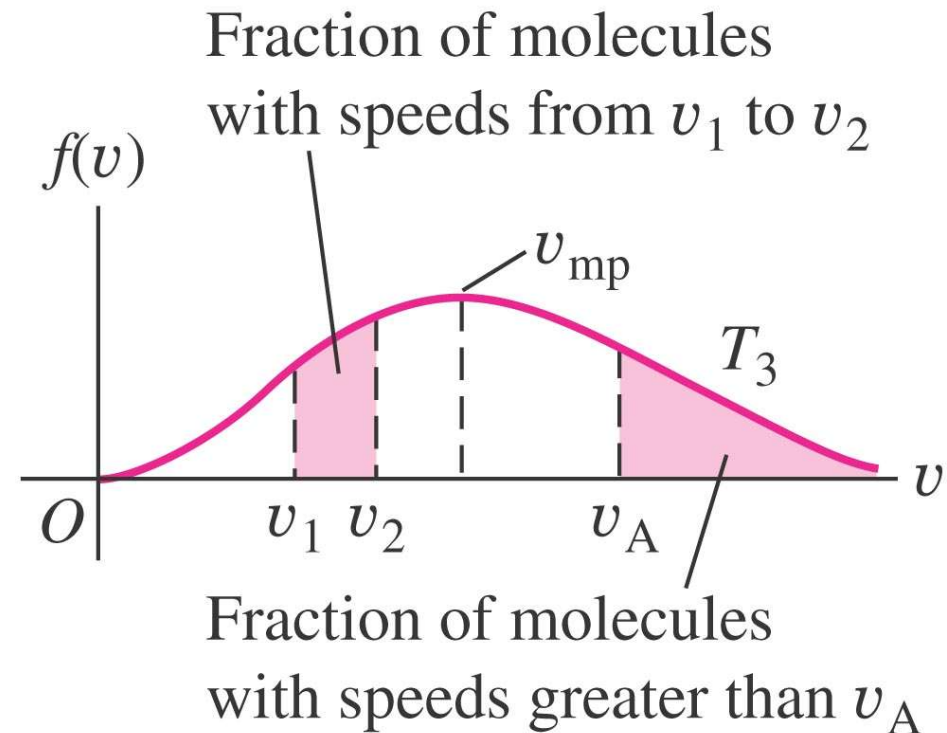
Histogram and Distribution

We all are familiar with the concept of using a histogram to represent a collection of different test scores.

Distribution of PHYS 262/266 Exam1 Scores

Distribution of Molecular Speeds

- Within an ideal gas, molecules moves with a diversity of speeds

- A mathematical rigorous way to describe this statistically is thru a **distribution function** $f(v)$.

- Properties:
  - $f(v) \sim 0$, $v \sim 0$
  - $f(v) \rightarrow 0$, $v \rightarrow large$
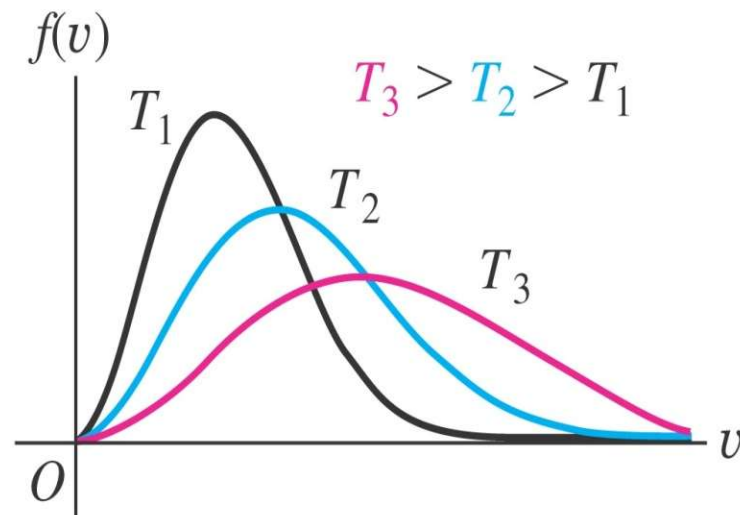 - $f(v)$ largest at mid-range
  - $f(v)$ fixed in time

Fraction of molecules with speeds from v_1 to v_2

$f(v)$

v_{mp}

T_3

O v_1 v_2 v_A v

Fraction of molecules with speeds greater than v_A

Maxwell-Boltzmann Distribution

$$f(\upsilon) = 4\pi \left(\frac{m}{2\pi kT} \right)^{3/2} \upsilon^2 e^{-m\upsilon^2/2kT}$$

□ $f(v)dv$ gives the probability of finding molecules with speed in range $[v, v+dv]$.

□ Diff averages with respect to the distribution of molecular speeds can be calculated using $f(v)$:

1. $v_{av} = \int_0^\infty vf(v)dv$ (avg of v)

2. $(v^2)_{av} = \int_0^\infty v^2 f(v)dv$ (avg of v^2)

$f(\upsilon)$

$T_3 > T_2 > T_1$

T_1

T_2

T_3

O v

As temperature increases:
• the curve flattens.
• the maximum shifts to higher speeds.

Example: Test Scores

4 Test Scores: $s_i = \{65, 80, 80, 95\}, i = 1, \cdots, 4$

Simple Avg: $S_{av} = \dfrac{1}{N}\sum\limits_{i=1}^{N} s_i = \dfrac{1}{4}(65 + 80 + 80 + 95)$
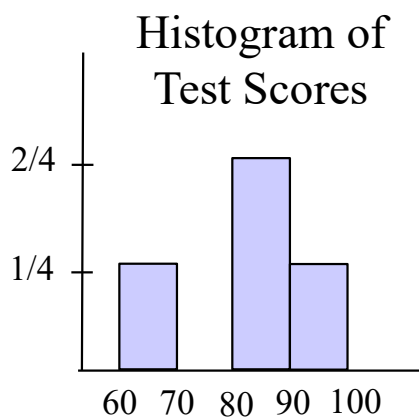
$S_{av} = \dfrac{1}{4}(65) + \dfrac{2}{4}(80) + \dfrac{1}{4}(95) = 80$ weighted average

Histogram of Test Scores

2/4

1/4

60 70 80 90 100

$S_{av} = \sum f_i s_i$ where f_i gives the fraction of occurrence of s_i

$v_{av} = \displaystyle\int_{0}^{\infty} f(v)\,v\,dv$ where $f(v)dv$ gives the fraction of molecules moving with speed $[v, v+dv]$

Example: Test Scores

4 Test Scores: $s_i = \{65, 80, 80, 95\}, i = 1, \cdots, 4$

Simple Avg: $\left(S^2\right)_{av} = \dfrac{1}{N}\displaystyle\sum_{i=1}^{N} s_i^2 = \dfrac{1}{4}\left(65^2 + 80^2 + 80^2 + 95^2\right)$

$$\left(S^2\right)_{av} = \dfrac{1}{4}\left(65^2\right) + \dfrac{2}{4}\left(80^2\right) + \dfrac{1}{4}\left(95^2\right) = 6512.5$$

$$\left(S^2\right)_{av} = \sum f_i s_i^2 \quad \neq \quad \left(S_{av}\right)^2 = 80^2 = 6400$$

$$(v^2)_{av} = \int_0^\infty f(v) v^2 dv \neq (v_{av})^2$$

Statistical Description of Molecular Speed

☐   Average speed (mean value): $v_{av} = \sqrt{\dfrac{8kT}{\pi m}}$

☐   Root Mean Square (RMS) speed: $v_{rms} = \sqrt{\left(v^2\right)_{av}} = \sqrt{\dfrac{3kT}{m}}$

   Note: v_{av} does not equal v_{rms} !

In addition to simple "mean" and "median", there are other ways to statistically describe the "average" values for a distribution of molecules moving at different speeds!

☐   Most Probable Speed - the maximum value of the distribution function $f(v)$: $v_{mp} = \sqrt{\dfrac{2kT}{m}}$

$D_{drinking}$

Heat Capacities of Gases (at constant V)

(a)

(b)
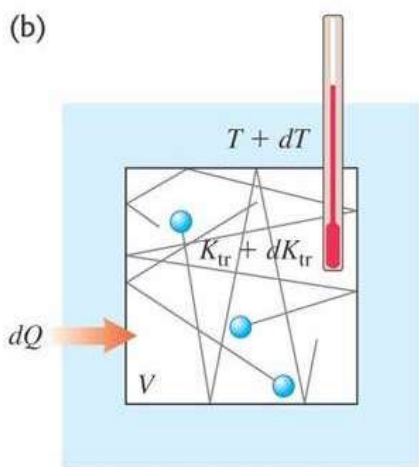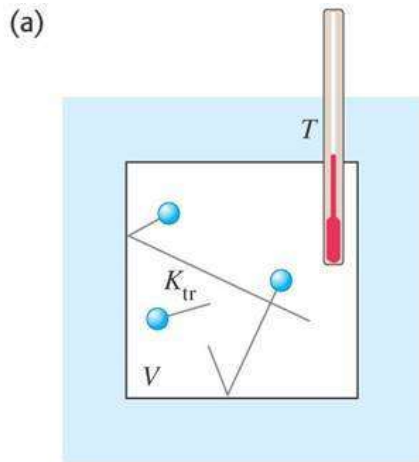
Using the Kinetic-Molecular model, one can calculate heat capacity for an *Ideal Gas*!

☐ For point-like molecules (monoatomic gases), molecular energy consists only of translational kinetic energy K_{tr}

☐ We just learned that K_{tr} is directly proportional to T.

☐ When an infinitesimal amount of heat dQ enters the gas, dT increases, and dK_{tr} increases accordingly,

$$dK_{tr} = \frac{3}{2}NkdT \quad (or = \frac{3}{2}nRdT)$$

Heat Capacities of Gases

□ From definition of molar heat capacity, we also know:

$$dQ \equiv nC_v dT$$

□ From mech-equivent of heat, requiring $dQ=dK_{tr}$ (*no work*) gives,

$$nC_v dT = \frac{3}{2}nRdT$$

□ *Monatomic* ~ Ideal Gas (matches well with prediction)

⟹ $C_v = \frac{3}{2}R$ $(=12.47J/mol \cdot K)$
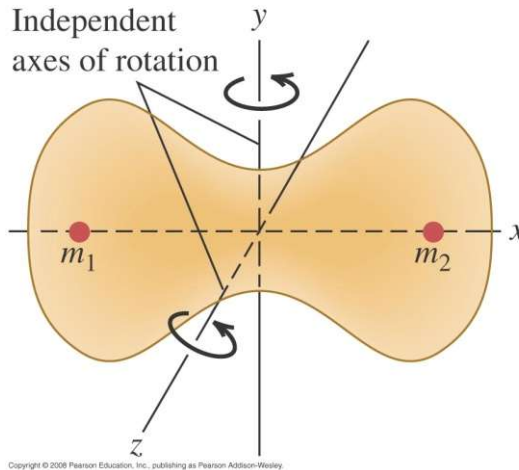
Table 18.1 Molar Heat Capacities of Gases

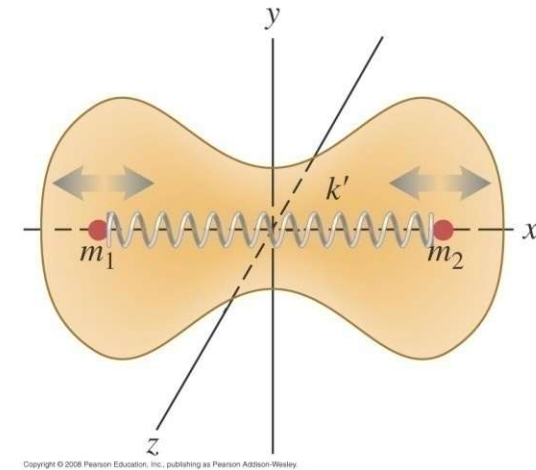| Type of Gas	Gas	C_V (J/mol · K)
Monatomic	He	12.47
	Ar	12.47
Diatomic	H_2	20.42
	N_2	20.76
	O_2	21.10
	CO	20.85
Polyatomic	CO_2	28.46
	SO_2	31.39
	H_2S	25.95

Heat Capacity (diatomic)

Translational motion. The molecule moves as a whole; its velocity may be described as the x-, y-, and z-velocity components of its center of mass.
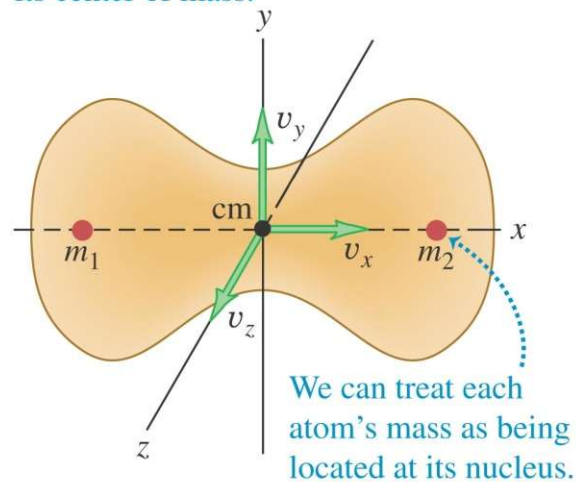
Rotational motion. The molecule rotates about its center of mass. This molecule has independent axes of rotation.

Vibrational motion. The molecule oscillates as though the nuclei were connected by a spring.

Independent axes of rotation

cm

m_1 m_2

v_y v_x v_z

We can treat each atom's mass as being located at its nucleus.

y x z

m_1 m_2

y x z

m_1 m_2 k'

y x z

A *diatomic* molecule can absorb energy into its translational motion, its rotational motion and in its vibrational motions.

Equipartition of Energy

This principle states that each **degree of freedom** ("separate mechanisms in storing energy") will contribute ($\frac{1}{2} kT$) to the total average energy per molecule.

☐ Monoatomic: 3 translational dofs → 3 ($\frac{1}{2} kT$)

This give $E_{tot} = 3/2\ NkT$ (same as before).

☐ Diatomic (*without* vibration): 3 trans dofs + 2 rotational dofs
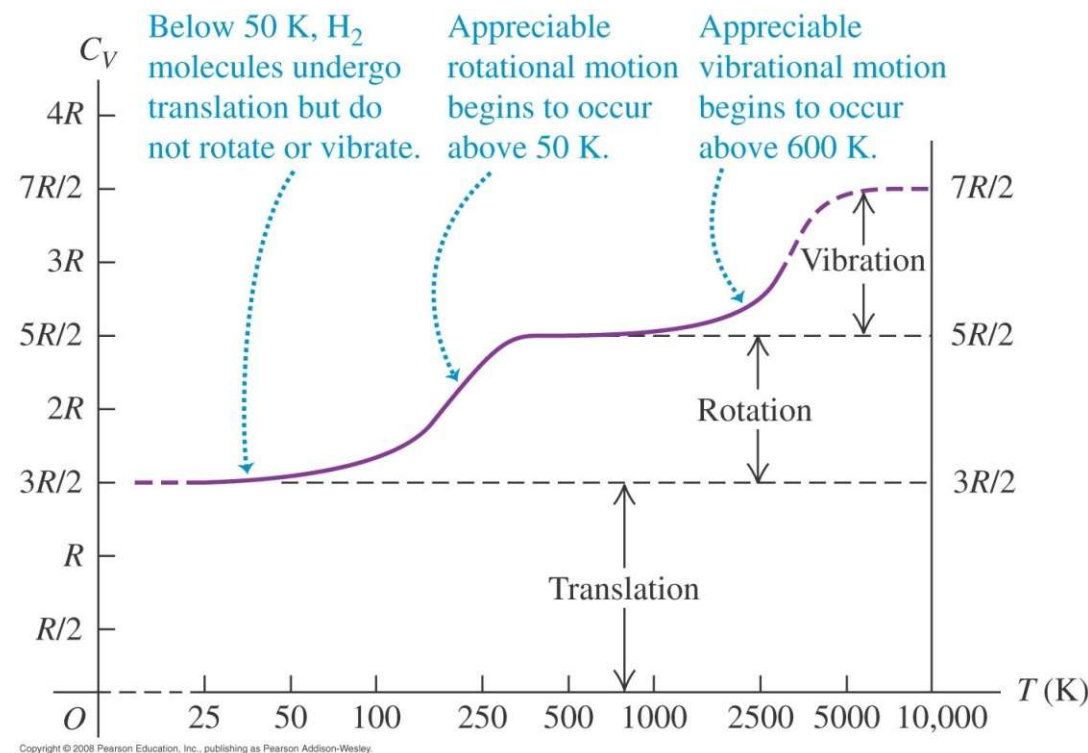
This give $E_{tot} = 5/2\ NkT$ or $= 5/2\ nRT$.

Again, consider an infinitesimal energy change, we have

$nC_v dT = \frac{5}{2} nRdT$, and this gives $C_v = 5/2\ R$.
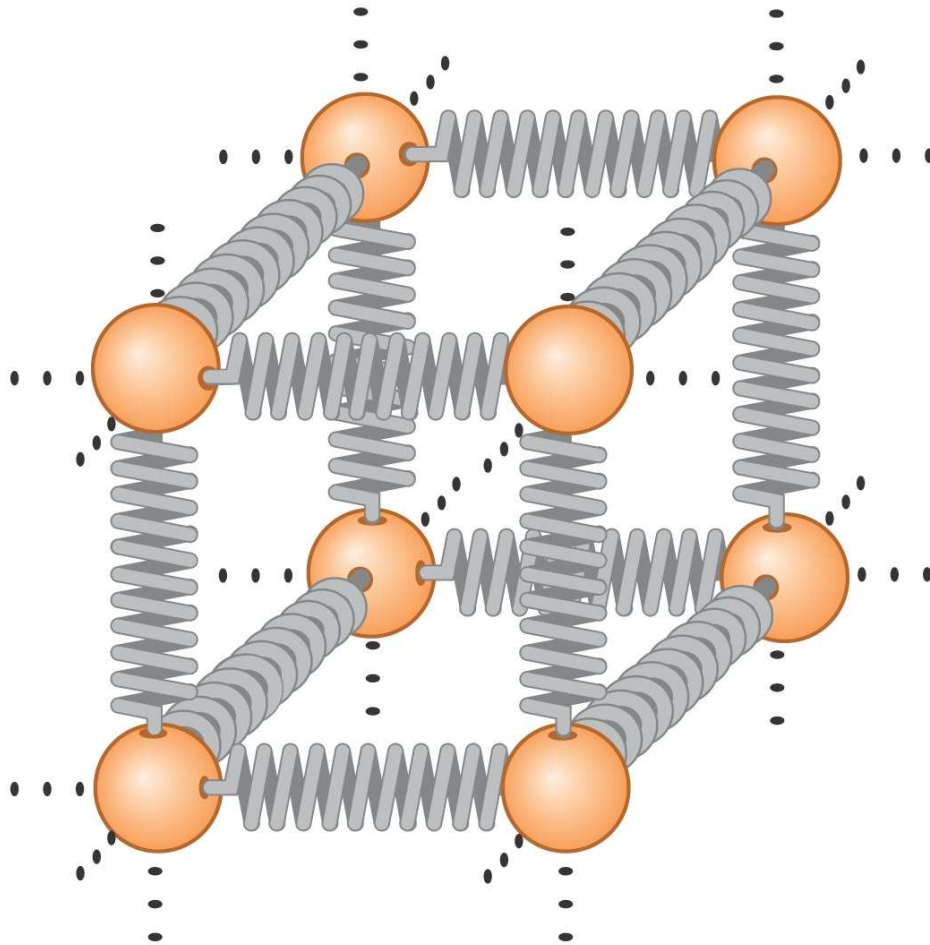
Heat Capacities (real gases, e.g., H$_2$)

- [ ] At low T, only 3 translational dofs can be activated

- [ ] At higher T, additional rotational dofs can be activated

- [ ] At higher T still, vibrational dofs might also get activated

Heat capacity for a H$_2$ gas

Below 50 K, H$_2$ molecules undergo translation but do not rotate or vibrate.

Appreciable rotational motion begins to occur above 50 K.

Appreciable vibrational motion begins to occur above 600 K.

C_V

4R
7R/2
3R
5R/2
2R
3R/2
R
R/2
O

Vibration

Rotation

Translation

7R/2
5R/2
3R/2

25   50   100   250   500   1000   2500   5000   10,000 T (K)

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→ For normal T range, one take $C_v = \dfrac{5}{2}R$ for H$_2$ gas
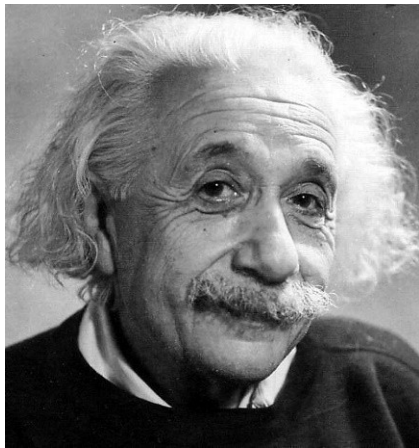
Heat Capacities of Ideal Solids

- ☐ Atoms are connected together by springs

- ☐ Assume harmonic motions for these springs (Hooke's law)

- ☐ For each spatial direction, there are two dofs (vibrational KE, vibrational PE) and we have 3 dims

- ☐ For the entire solid, we have 3/2 kT (KE) + 3/2 kT (PE) → $E_{tot} = 3kT$
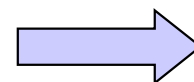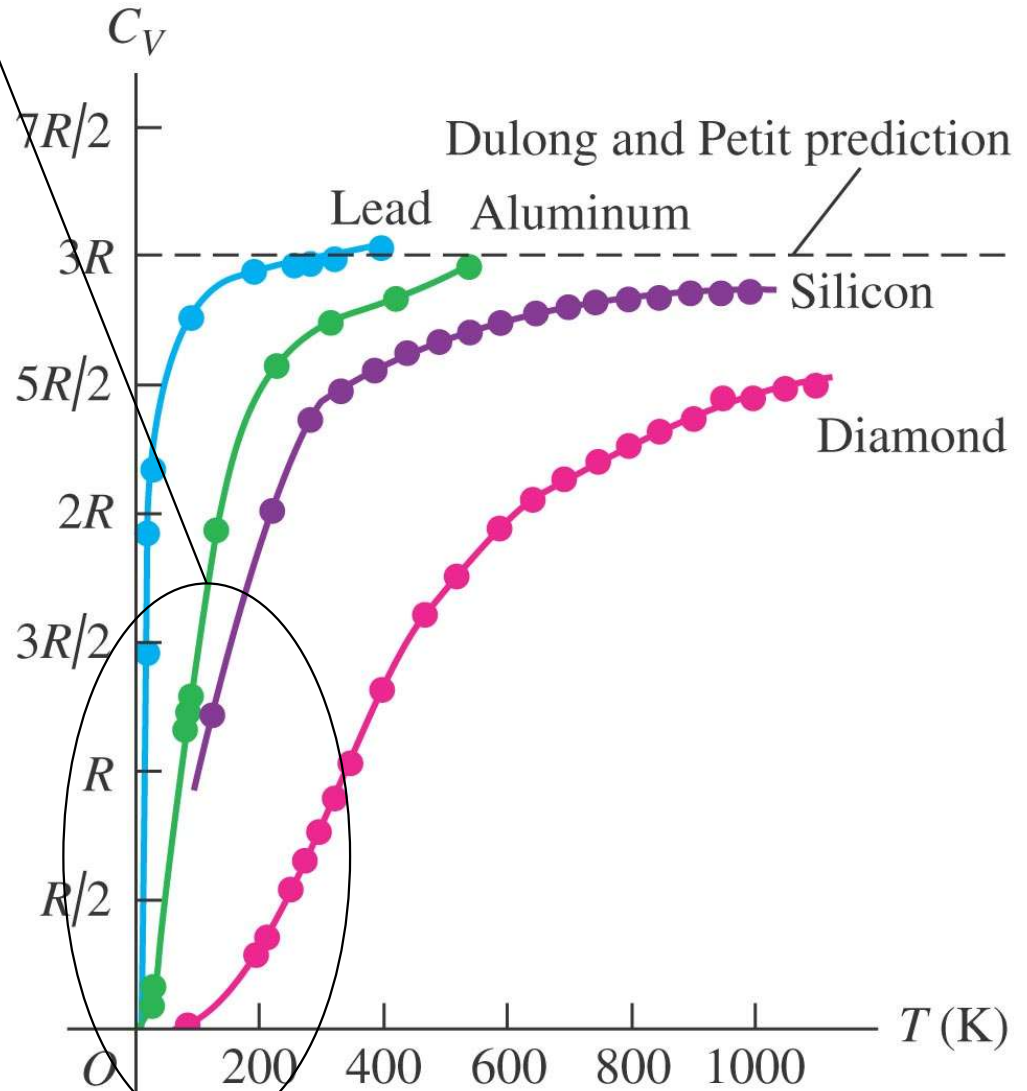
$$C_v = 3R$$

Dulong Petit Prediction and Improvements

Einstein: treating atoms as QM HMOs…

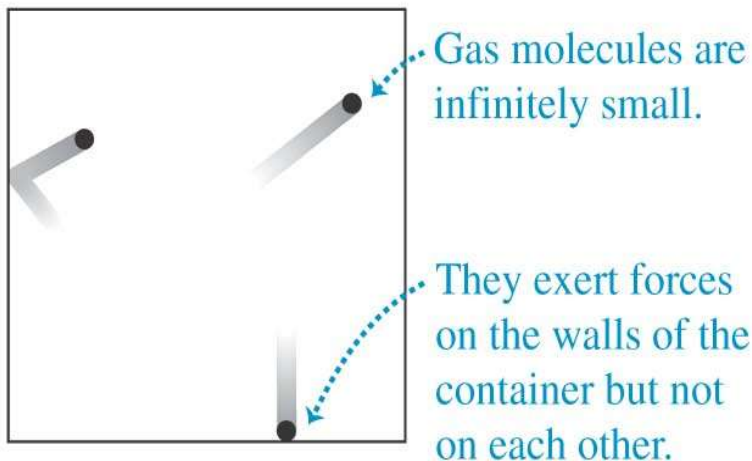$$C_v \to 0$$

Debye: extension by including low-f phonons…

$$C_v \sim T^3$$

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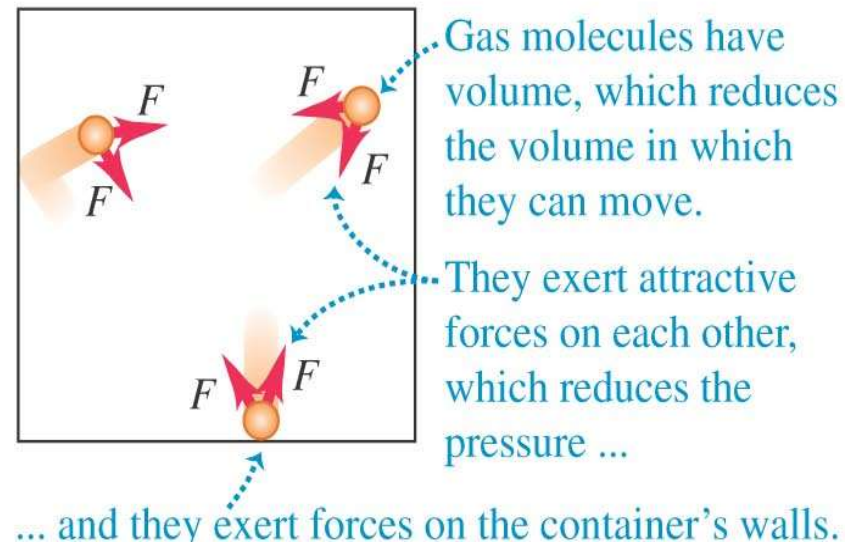
The van der Waals Equation

A more realistic Equation of States for gases which includes corrections for the facts that molecules are not point particles, that they have volume, and for the attraction/repulsion that naturally exists between the adjacent atoms/molecules.

(a) An idealized model of a gas

Gas molecules are infinitely small.

They exert forces on the walls of the container but not on each other.

(b) A more realistic model of a gas

F F F F F F

Gas molecules have volume, which reduces the volume in which they can move.

They exert attractive forces on each other, which reduces the pressure ...

... and they exert forces on the container's walls.

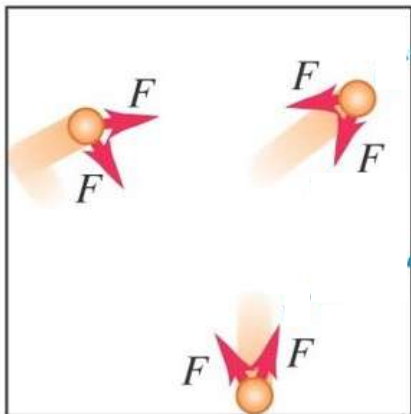$$\left(P+\frac{an^2}{V^2}\right)(V-bn)=nRT$$

Notes on Van Der Waals

- The volume parameter b:
  - It makes sense that real gas as finite size hard spheres will reduce the total volume of the gas by a term which is proportional to the number of mole n.
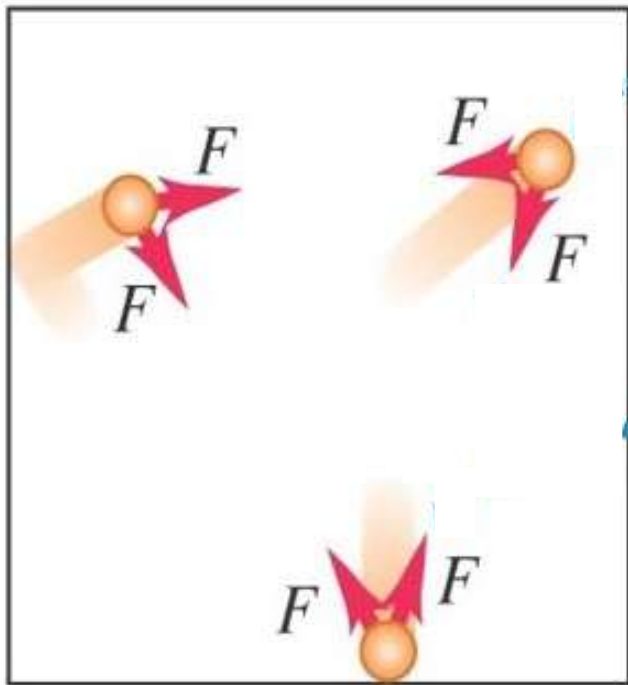
$$P = \frac{nRT}{V_{eff}} = \frac{nRT}{(V - bn)}$$

Notes on Van Der Waals

□ The intra-molecular force parameter a:

  ■ Intra-molecular force tends to *reduce* the pressure of the gas onto the wall by pulling the molecules toward the interior of the container

  ■ This intra-molecular force acts in pairs (to the lowest order of approximation)

  → depends on the number of *pairs* of molecules within this unit volume

  ■ For N molecules, # pairs = $N(N-1)/2$; for N large, $\sim N^2$
  The count of molecular pairs (in mole) within an unit volume $\sim (n/V)^2$

$$P = \frac{nRT}{(V - nb)} - a\left(\frac{n}{V}\right)^2$$

$$\rightarrow \quad \left(P + a\left(\frac{n}{V}\right)^2\right)(V - nb) = nRT$$